#### **Fourier Series Analysis**

MODIFIED BY TLH

**Trig and Exponential Series** 



Joseph Fourier lived from 1768 to 1830

**Fourier** studied the mathematical theory of heat conduction. He established the partial differential equation governing heat diffusion and solved it by using infinite series of trigonometric functions.

On the interval [-T/2, T/2], the limits of integration for the Fourier can be changed from  $[-\pi, \pi]$  by assigning to the integration variety value  $2\pi t/T$ . The period of the function is thus T.

Assuming that f(t) is continuous on the interval  $-T/2 \le t \le t$  the coefficients  $a_n$  and  $b_n$  can be computed by the formulas

 $a_0 = \frac{2}{T} \int_{-T/2}^{T/2} f(t) dt,$ 

$$a_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos\left(\frac{2n\pi t}{T}\right) dt,$$

$$b_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin\left(\frac{2n\pi t}{T}\right) dt,$$

where n = 1, 2, ... is any positive integer.

The Fourier series on the interval [-T/2, T/2] is thus

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[ a_n \cos\left(\frac{2n\pi t}{T}\right) + b_n \sin\left(\frac{2n\pi t}{T}\right) \right].$$

#### $\square$ EXAMPLE 8.4 Fourier series square wave example

A square wave of amplitude A and period T shown in Figure 8.4 can be defined as

$$f(t) = \begin{cases} A, & 0 < t < \frac{T}{2}, \\ -A, & -\frac{T}{2} < t < 0, \end{cases}$$

with f(t) = f(t+T), since the function is periodic.

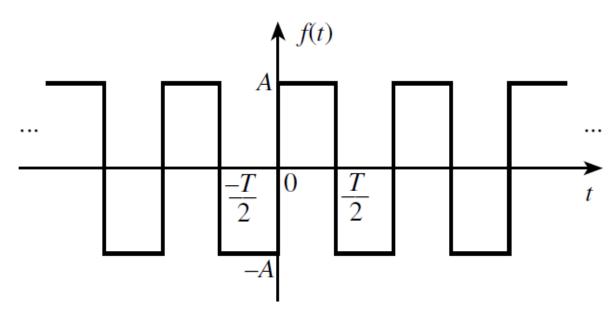


FIGURE 8.4 Square wave of Example 8.4

The first observation is that f(t) is odd, which yields the result that  $a_0 = 0$  and  $a_i = 0$  for every coefficient of the cosine terms. Letting  $\omega_0 = 2\pi/T$ , the coefficients  $b_n$  are

$$b_n = 2\left(\frac{2}{T}\right) \int_0^{T/2} A \sin(n\omega_0 t) dt.$$

The result is

$$f(t) = \frac{4A}{\pi} \sum_{n=1}^{\infty} \frac{\sin[(2n-1)\omega_0 t]}{(2n-1)},$$

where (2n-1) is introduced to assure that only odd terms are included in the summation. The sine waves that make up the Fourier series for the odd square wave are

$$f(t) = \frac{4A}{\pi} \left[ \sin(\omega_0 t) + \frac{\sin(3\omega_0 t)}{3} + \cdots \right],$$

## Harmonic Signal->Periodic

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi k F_0 t}$$

Sums of <u>Harmonic</u> complex exponentials are <u>Periodic</u> signals

#### PERIOD/FREQUENCY of COMPLEX EXPONENTIAL:

$$2\pi(F_0) = \omega_0 = \frac{2\pi}{T_0}$$
 or  $T_0 = \frac{1}{F_0}$ 

#### ☐ EXAMPLE 8.5

#### Complex Series Square Wave Example

Consider the odd square wave of Example 8.4 and the complex Fourier coefficients

$$\alpha_n = \frac{1}{T} \int_{-T/2}^{0} (-A)e^{-in\omega_0 t} dt + \frac{1}{T} \int_{0}^{T/2} (A)e^{-in\omega_0 t} dt, \qquad (8.29)$$

which leads to the series

$$f(t) = \frac{2A}{i\pi} \sum_{n=-\infty}^{\infty} \frac{e^{i(2n-1)\omega_0 t}}{(2n-1)},$$
(8.30)

as defined in Equation 8.23.

This form contains complex coefficients, but the series can be written in terms of sine waves by combining the corresponding terms for positive and negative arguments. To determine the coefficients, the amount of difficulty is about the same for the trigonometric series and the complex series. However, the complex series perhaps has an advantage when the magnitude of the coefficients are of interest.

Each coefficient has the form

$$\alpha_n = \frac{2A}{in\pi} = \frac{2A}{n\pi}e^{-i\pi/2}, \qquad n = \pm 1, \pm 3, \dots,$$

and the coefficients for even values,  $n=0,\pm 2,\ldots,$  are zero. Notice that the

The trigonometric series is derived from the complex series by expanding the complex series of Equation 8.30 as

$$f(t) = \sum_{n=-\infty}^{\infty} \alpha_n e^{in\omega_0 t}$$

$$= \cdots - \frac{2A}{3\pi i} e^{-i3\omega_0 t} - \frac{2A}{\pi i} e^{-i\omega_0 t} + \frac{2A}{\pi i} e^{i\omega_0 t} + \frac{2A}{3\pi i} e^{i3\omega_0 t} + \cdots$$

and recognizing the sum of negative and positive terms for each n as  $2\sin(n\omega_0 t)$ . The trigonometric series becomes

$$f(t) = \frac{4A}{\pi} \left( \sin(\omega_0 t) + \frac{\sin(3\omega_0 t)}{3} + \cdots \right) = \frac{4A}{\pi} \sum_{n=1}^{\infty} \frac{\sin[(2n-1)\omega_0 t]}{(2n-1)},$$

which is the result of Example 8.4.

#### LECTURE OBJECTIVES

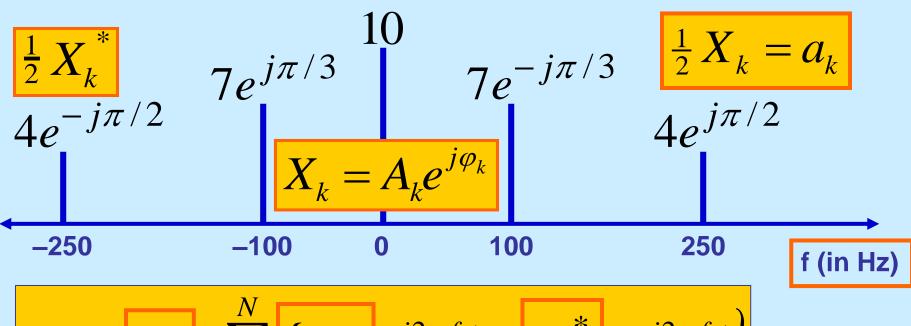
Work with the Fourier Series Integral

$$a_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-j(2\pi k/T_0)t} dt$$

- ANALYSIS via Fourier Series
  - For <u>PERIODIC</u> signals:  $x(t+T_0) = x(t)$
  - Draw spectrum from the Fourier Series coeffs

#### SPECTRUM DIAGRAM

Recall Complex Amplitude vs. Freq



$$x(t) = a_0 + \sum_{k=1}^{N} \left\{ a_k e^{j2\pi f_k t} + a_k^* e^{-j2\pi f_k t} \right\}$$

## **Harmonic Signal->Periodic**

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi k F_0 t}$$

Sums of <u>Harmonic</u> complex exponentials are <u>Periodic</u> signals

#### PERIOD/FREQUENCY of COMPLEX EXPONENTIAL:

$$2\pi(F_0) = \omega_0 = \frac{2\pi}{T_0}$$
 or  $T_0 = \frac{1}{F_0}$ 

# Notation for Fundamental Frequency in Fourier Series

- The k-th frequency is  $f_k = kF_0$ 

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi k F_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi f_k t}$$

- Thus,  $f_0 = 0$  is DC
- This is why we use upper case F<sub>0</sub> for the Fundamental Frequency

# **STRATEGY:** $x(t) \rightarrow a_k$

#### ANALYSIS

- Get representation from the signal
- Works for <u>PERIODIC</u> Signals
- Measure similarity between signal & harmonic
- Fourier Series
  - Answer is: an INTEGRAL over one period

$$a_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-j\omega_0 kt} dt$$

#### **CALCULUS for complex exp**

$$\frac{d}{dt}e^{\alpha t} = \alpha e^{\alpha t} \quad \to \quad \frac{d}{dt}e^{j\alpha t} = j\alpha e^{\alpha t}$$

$$\int_{a}^{b} e^{\beta t} dt = \frac{1}{\beta} e^{\beta t} \bigg|_{a}^{b} = \frac{1}{\beta} \left( e^{\beta b} - e^{\beta a} \right)$$

$$\int_{a}^{b} e^{j\beta t} dt = \frac{1}{j\beta} e^{j\beta t} \bigg|_{a}^{b} = \frac{1}{j\beta} \left( e^{j\beta b} - e^{j\beta a} \right)$$

## **Fourier Series Integral**

Use orthogonality to determine a<sub>k</sub> from x(t)

$$a_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-j(2\pi/T_0)kt} dt$$

Fundamental Freq.

$$F_0 = 1/T_0$$

$$a_{-k} = a_k^*$$
 when  $x(t)$  is real

$$a_0 = \frac{1}{T_0} \int_{0}^{T_0} x(t) dt$$

THIS IS THE AVERAGE!

(DC component)

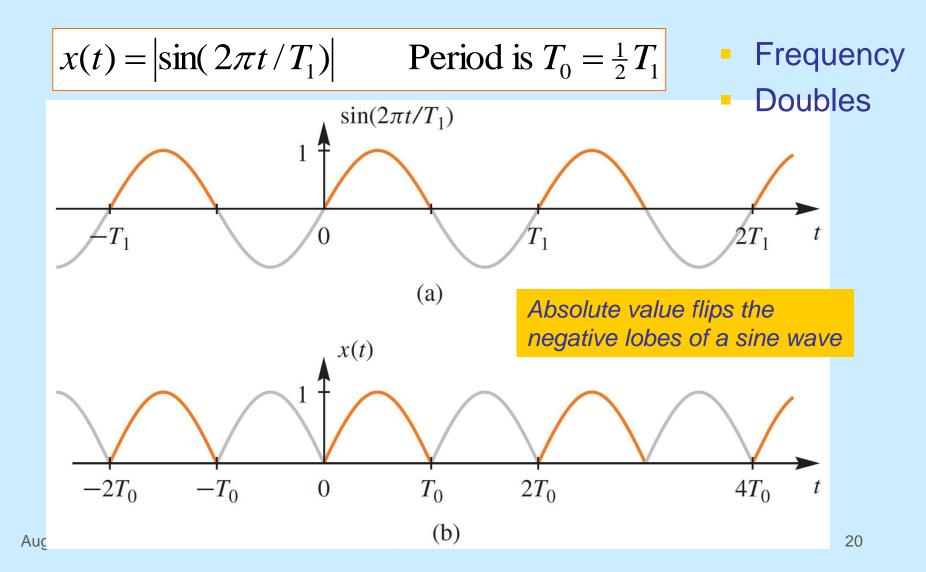
# Fourier Series: $x(t) \rightarrow a_k$

- ANALYSIS
  - Given a <u>PERIODIC Signal</u>
  - Fourier Series coefficients are obtained via an <u>INTEGRAL over one period</u>

$$a_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-j\omega_0 kt} dt$$

- Next, consider a specific signal, the FWRS
  - Full Wave Rectified Sine

#### **Full-Wave Rectified Sine**



## Full-Wave Rectified Sine {a<sub>k</sub>}

$$a_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-j(2\pi/T_0)kt} dt$$

$$a_{k} = \frac{1}{T_{0}} \int_{0}^{T_{0}} \sin(\frac{\pi}{T_{0}}t) e^{-j(2\pi/T_{0})kt} dt$$

$$= \frac{1}{T_{0}} \int_{0}^{T_{0}} \frac{e^{j(\pi/T_{0})t} - e^{-j(\pi/T_{0})t}}{2j} e^{-j(2\pi/T_{0})kt} dt$$

$$= \frac{1}{j2T_0} \int_{0}^{T_0} e^{-j(\pi/T_0)(2k-1)t} dt - \frac{1}{j2T_0} \int_{0}^{T_0} e^{-j(\pi/T_0)(2k+1)t} dt$$

$$= \frac{e^{-j(\pi/T_0)(2k-1)t}}{j2T_0(-j(\pi/T_0)(2k-1))} \Big|_0^{T_0} - \frac{e^{-j(\pi/T_0)(2k+1)t}}{j2T_0(-j(\pi/T_0)(2k+1))} \Big|_0^{T_0}$$

#### Full-Wave Rectified Sine

$$x(t) = \left| \sin(2\pi t / T_1) \right|$$

$$Period: T_0 = \frac{1}{2}T_1$$

$$\Rightarrow x(t) = \left| \sin(\pi t / T_0) \right|$$

# Full-Wave Rectified Sine {a<sub>k</sub>}

$$a_k = \frac{e^{-j(\pi/T_0)(2k-1)t}}{j2T_0(-j(\pi/T_0)(2k-1))} \Big|_0^{T_0} - \frac{e^{-j(\pi/T_0)(2k+1)t}}{j2T_0(-j(\pi/T_0)(2k+1))} \Big|_0^{T_0}$$

$$= \frac{1}{2\pi(2k-1)} \left( e^{-j(\pi/T_0)(2k-1)T_0} - 1 \right) - \frac{1}{2\pi(2k+1)} \left( e^{-j(\pi/T_0)(2k+1)T_0} - 1 \right)$$

$$= \frac{1}{\pi(2k-1)} \left( e^{-j\pi(2k-1)} - 1 \right) - \frac{1}{\pi(2k+1)} \left( e^{-j\pi(2k+1)} - 1 \right)$$

$$= \left(\frac{2k+1-(2k-1)}{\pi(4k^2-1)}\right)\left(-(-1)^{2k}-1\right) = \frac{-2}{\pi(4k^2-1)}$$

## Fourier Coefficients: ak

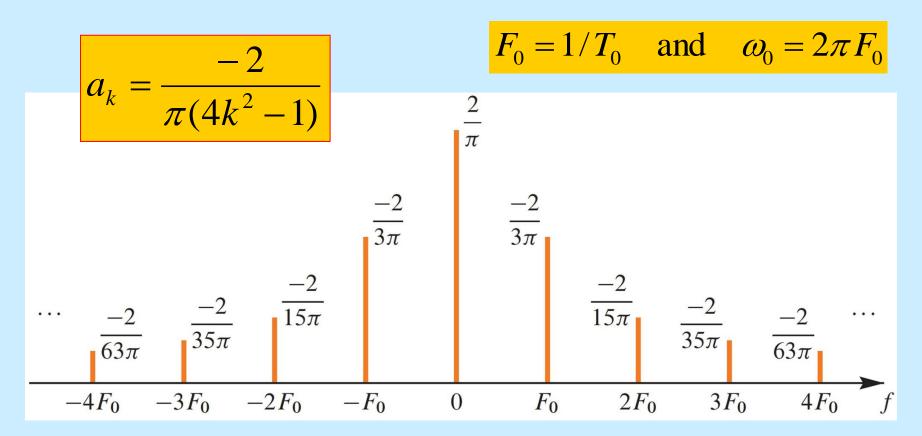
- a<sub>k</sub> is a function of k
  - Complex Amplitude for k-th Harmonic

$$a_k = \frac{-2}{\pi (4k^2 - 1)}$$
 NOTE:  $\frac{1}{k^2}$  for large  $k$ 

- Does not depend on the period, T<sub>0</sub>
- **DC** value is  $a_0 = 2/\pi = 0.6336$

## **Spectrum from Fourier Series**

Plot  $a_k$  for Full-Wave Rectified Sinusoid



## **Reconstruct From Finite Number** of Harmonic Components

Full-Wave Rectified Sinusoid  $x(t) = |\sin(\pi t/T_0)|$ 

$$T_0 = 10 \text{ ms}$$
  
 $\Rightarrow F_0 = 100 \text{ Hz}$ 

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$$x_N(t) = a_0 + \sum_{k=1}^{N} \left\{ a_k e^{j2\pi k F_0 t} + a_k^* e^{-j2\pi k F_0 t} \right\}$$

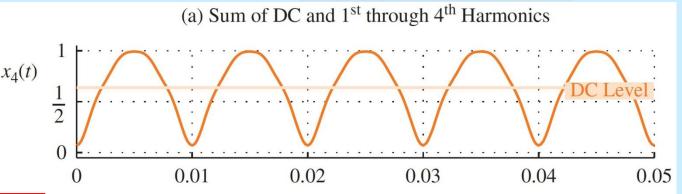
How close is 
$$x_N(t)$$
 to  $x(t) = \left| \sin(\pi t / T_0) \right|$ ?

## **Reconstruct From Finite Number** of Spectrum Components

Full-Wave Rectified Sinusoid  $x(t) = |\sin(\pi t/T_0)|$ 

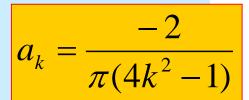
$$x(t) = \left| \sin(\pi t / T_0) \right|$$

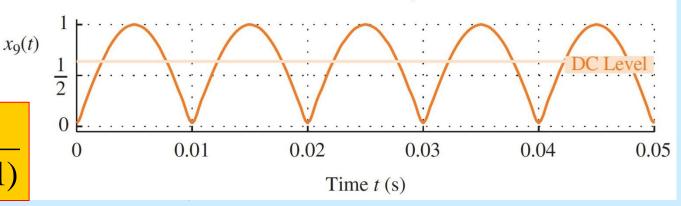
$$T_0 = 10 \text{ ms}$$
  
 $\Rightarrow F_0 = 100 \text{ Hz}$   $x_4(t)$ 



$$a_0 = 2/\pi = 0.6336$$

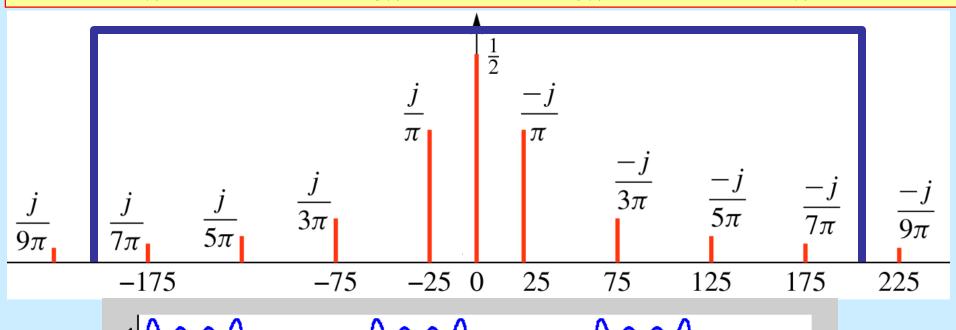
(b) Sum of DC and 1st through 9th Harmonics

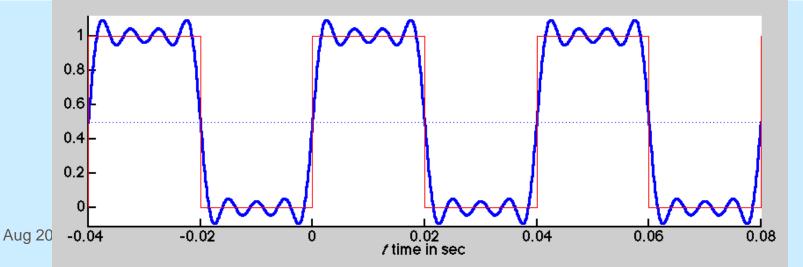




#### Synthesis: up to 7th Harmonic

$$y(t) = \frac{1}{2} + \frac{2}{\pi}\cos(50\pi t - \frac{\pi}{2}) + \frac{2}{3\pi}\sin(150\pi t) + \frac{2}{5\pi}\sin(250\pi t) + \frac{2}{7\pi}\sin(350\pi t)$$





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